

# MAXWELL'S EQUATIONS

ME in INTEGRAL FORM

Line integral

Surface integral

ME in DIFFERENTIAL FORM

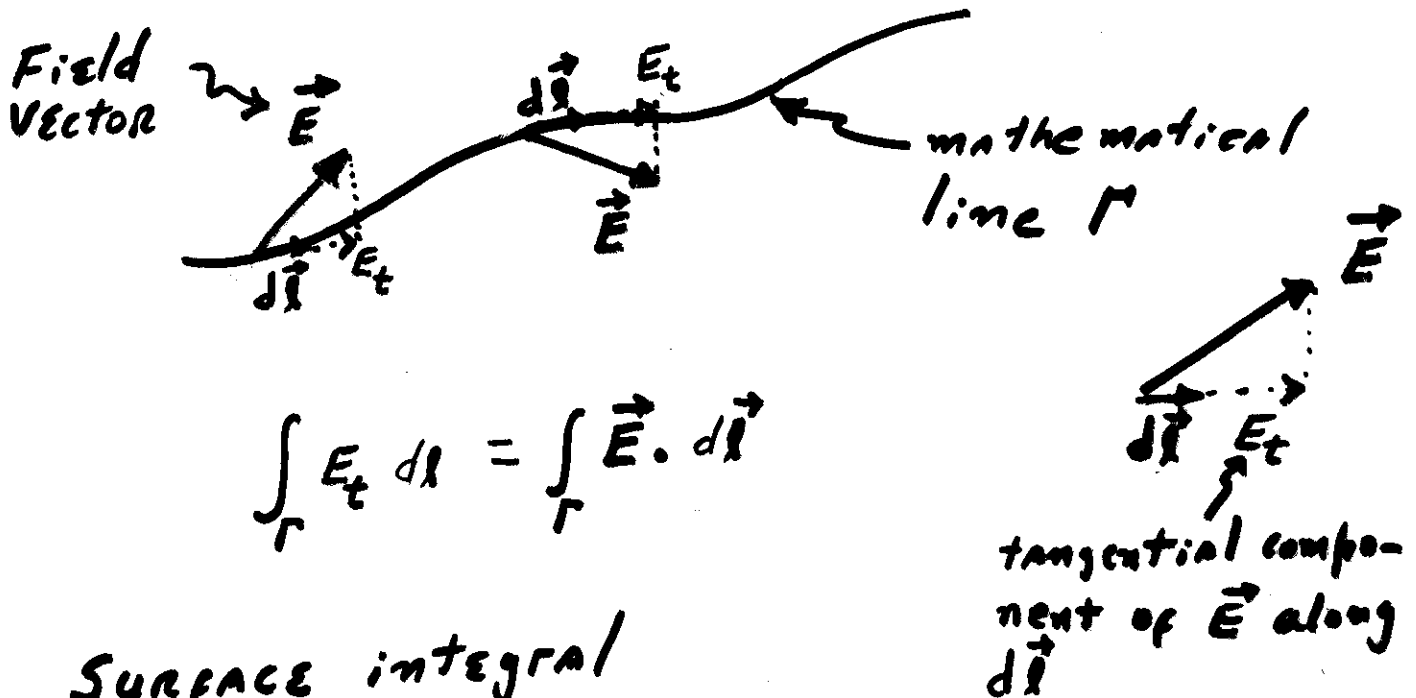
Operators GRADIENT  
DIVERGENCE  
ROTATIONAL

GAUSS' theorem

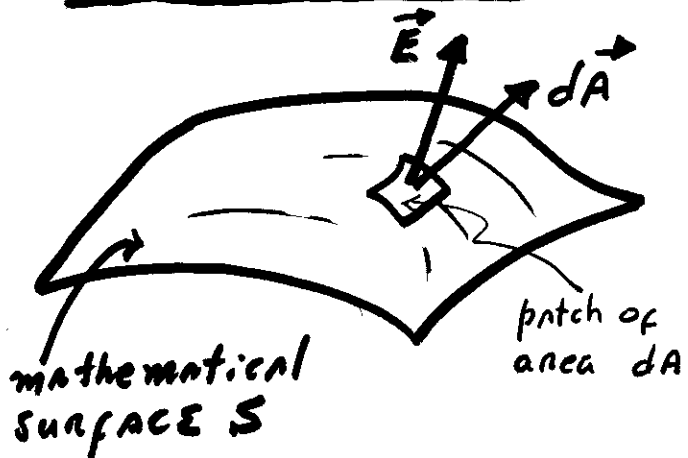
Stoke's theorem

GENERATION, PROPAGATION and  
DETECTION of ELECTROMAGNETIC WAVES

## Line integral



## SURFACE integral

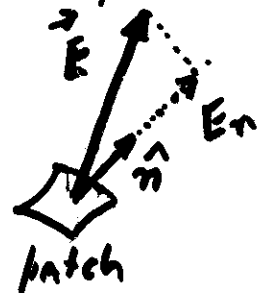


$$d\vec{A} = \hat{n} dA$$

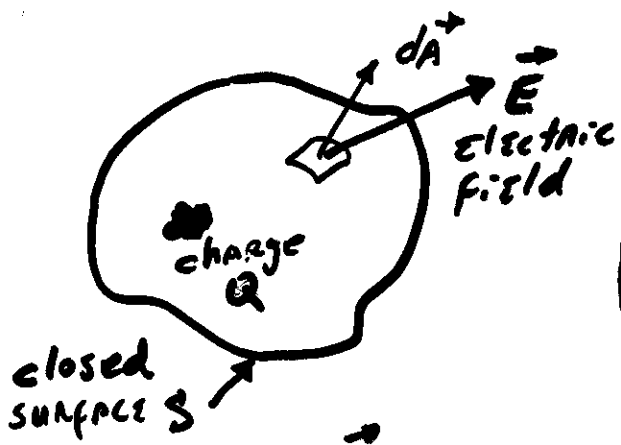
where

$\hat{n}$  is the unit vector perpendicular to the patch

$$\int_S E_n dA = \int_S \vec{E} \cdot \hat{n} dA = \int_S \vec{E} \cdot d\vec{A}$$

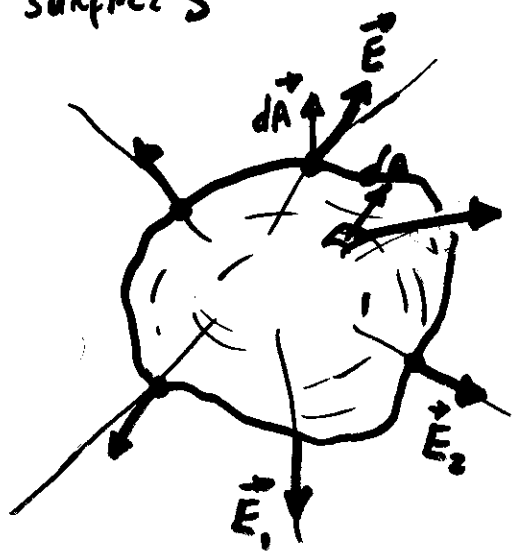


First Maxwell Eq



$$\int_S \vec{E} \cdot d\vec{A} = \frac{Q_{\text{inside}}}{\epsilon_0}$$

$\epsilon_0$ : permittivity of free-space



$\vec{E}$  is evaluated at each point ON the surface  $S$

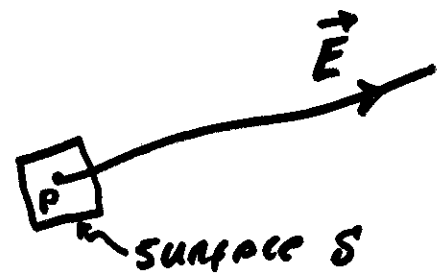
Electric flux  $\phi_E = \int_S \vec{E} \cdot d\vec{A}$

Notice:



If an electric field line originates at a given point P, then there must exist charge at that point

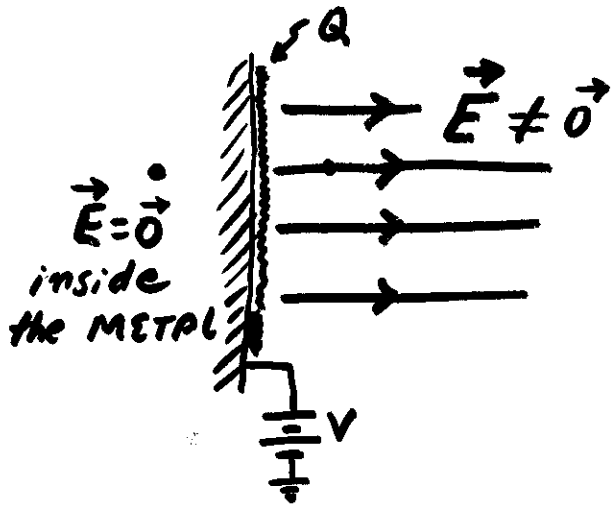
Proof



obviously  $\phi_E \neq 0$

Therefore, according to the 1st ME, there must be charge at P

This is what happens in metals, for example. 4



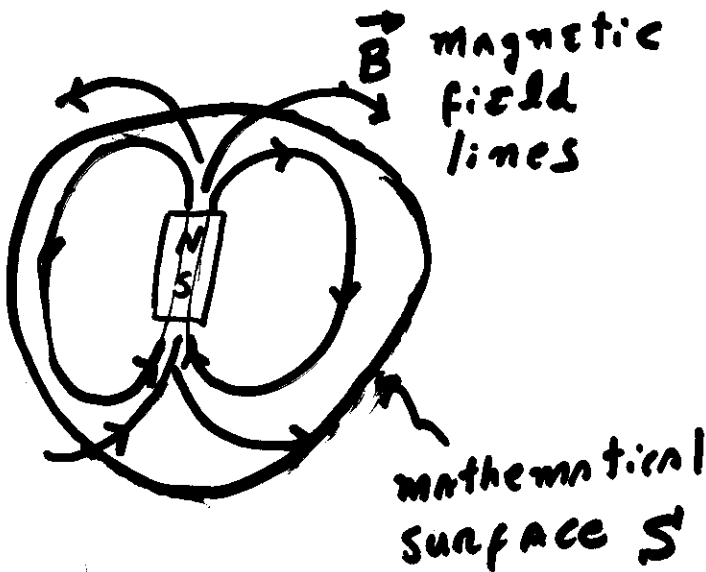
When charges are present:

lines of  $\vec{E}$  originate on positive charges and terminate on negative charges

Everywhere else

the  $\vec{E}$  lines can twist and turn in space, but they cannot start or stop

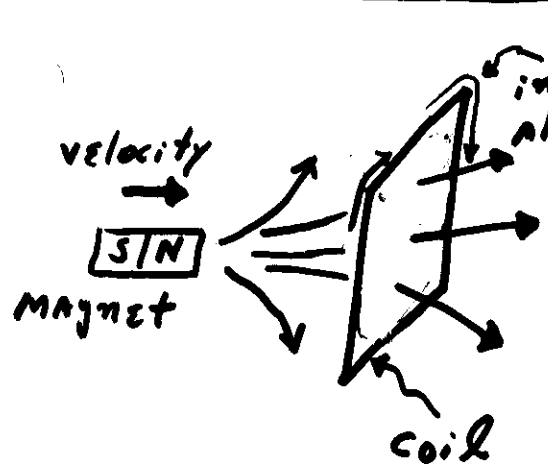
## SECOND MAXWELL Eq.



$$\int_S \vec{B} \cdot d\vec{A} = 0$$

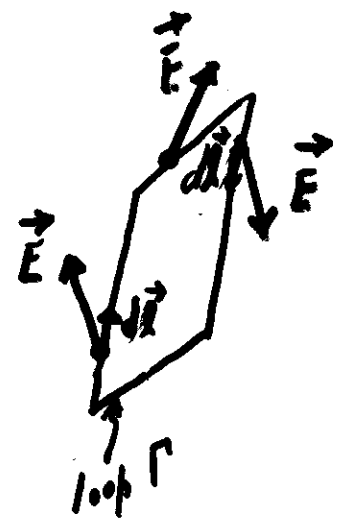
No magnetic monopoles have been observed (so far)

# Third Maxwell Eq



The motion of the magnet produces an induced current along the coil

The existence of a current implies the presence of a electric field or (equivalently) an electromotive force



$$\mathcal{E} = \int_{\text{loop } \Gamma} \vec{E} \cdot d\vec{l} \quad (= Ri)$$

electromotive force

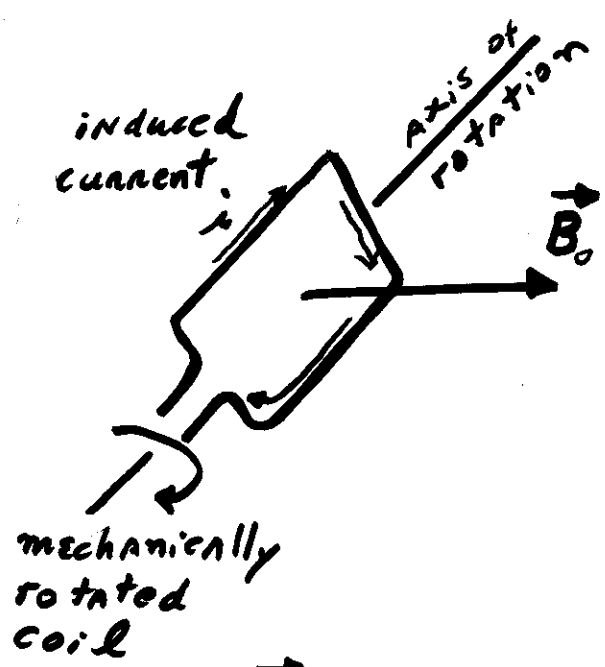
But, what is the value of  $\mathcal{E}$ ?  
(1V, 2V, -2.3V...?)

$$\mathcal{E} = \int_{\text{loop } \Gamma} \vec{E} \cdot d\vec{l} = -\frac{d}{dt} \left( \int_S \vec{B} \cdot d\vec{A} \right)$$

where  $S$  is any open surface having the loop  $\Gamma$  as its boundary

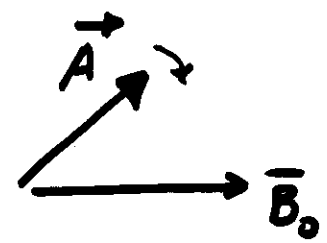
$$\mathcal{E} = -\frac{d}{dt} \Phi_M$$





Notice:  
 A  $\vec{B} = \vec{B}(t)$  ensures the induction of an electric field  $\vec{E}$  and, consequently, a current  $i$  along a coil.

However, a static magnetic field  $\vec{B}_0$  can also be exploited to generate currents along a coil (see figure).



$\vec{A}$  revolves around  $\vec{B}_0$  causing a time dependent magnetic flux  $\Phi_m$ .

So, to have an electromotive force what matters is to have time dependent  $\Phi_m(t)$

[In the figure above, since there is not an induced electric field, how can we have an electromotive force  $\mathcal{E} = \int \vec{E} \cdot d\vec{s}$ ? Shouldn't  $\vec{E}$  be zero because  $\vec{E} = \vec{0}$ ?

Answer: We better re-define  $\mathcal{E}$  as

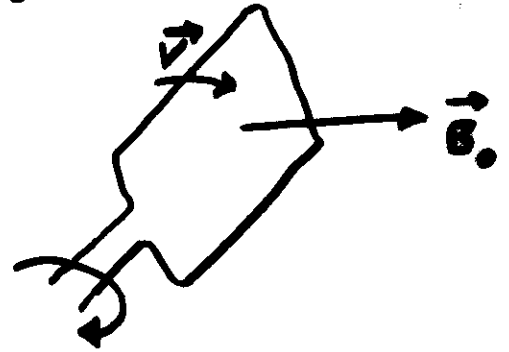
$$\mathcal{E} = \int_{\text{loop } \Gamma} \frac{\text{Magnetic Electric force } \vec{F} \text{ on } q}{q} \cdot d\vec{l}$$

For the particular case of a mechanically rotated coil immersed in a uniform magnetic field  $\vec{B}_0$ . (see previous figure)

$$\mathcal{E} = \int_{\text{loop } \Gamma} \frac{\text{magnetic force on } q}{q} \cdot d\vec{l}^{\rightarrow}$$

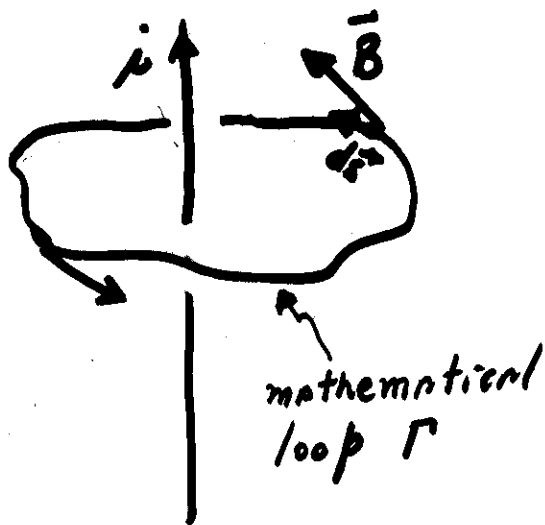
$$= \int_{\text{loop } \Gamma} \frac{q \vec{v} \times \vec{B}_0}{q} \cdot d\vec{l}^{\rightarrow}$$

$$= \int_{\text{loop } \Gamma} (\vec{v} \times \vec{B}_0) \cdot d\vec{l}^{\rightarrow} \quad ]$$



### 4th Maxwell's Eq.

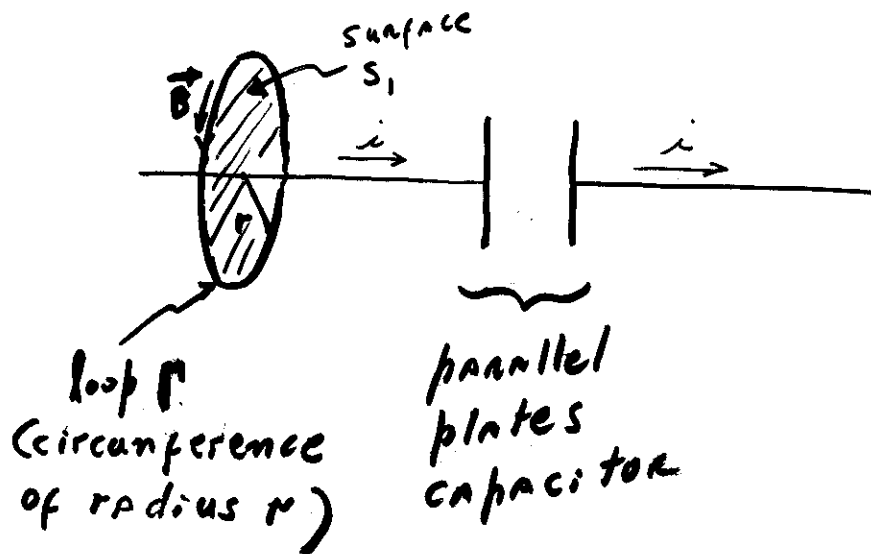
It built upon Ampere's Law



$$\int_{\text{loop } \Gamma} \vec{B} \cdot d\vec{s} = \mu_0 \underbrace{i}_{\text{current enclosed by the loop } \Gamma}$$

Maxwell notice this law was incomplete

James Clerk Maxwell noticed there was something wrong with the fourth equation. For example:

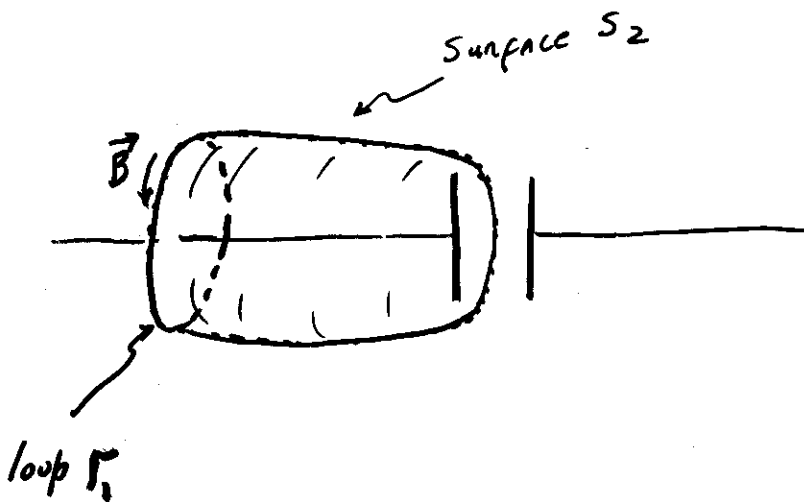


$$\int_{\Gamma} \vec{B} \cdot d\vec{l} = \mu_0 \underbrace{i}_{\text{current crossing the surface } S_1}$$

$$B 2\pi r = \mu_0 i$$

$$\Rightarrow \text{So, } B = \frac{\mu_0}{2\pi r} i \quad (1)$$

On the other hand,



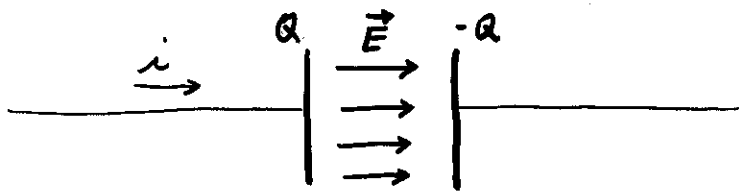
$$\int_{\Gamma} \vec{B} \cdot d\vec{l} = \mu_0 \underbrace{0}_{\text{current crossing the surface } S_2}$$

$$B 2\pi r = \mu_0 \times 0 \quad (2)$$

$$\text{So, } B = 0$$

How can be possible that (1) and (2) give different results for the same  $B$ ?  
Something must be wrong with the Ampere's law

In order to solve this contradictory situation, let's take a look to what is going on inside the parallel plates capacitor

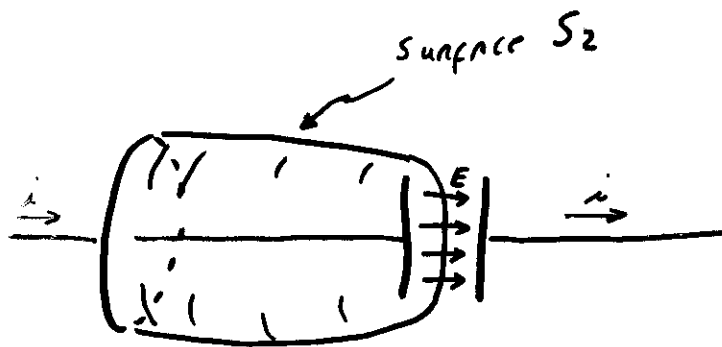


Let's remember

$$E = \frac{Q}{A \epsilon_0}$$

(A is the area of the plates)

The current  $i$  continuously accumulates charge on the plates, this is  $Q = Q(t)$



Notice, the electric flux  $\Phi_E$  crossing the surface  $S_2$  is given by

$$\Phi_E = EA = \frac{Q}{\epsilon_0}$$

or

$$Q = \epsilon_0 \Phi_E$$

From the last expression, we obtain

$$\frac{dQ}{dt} = \epsilon_0 \frac{d\Phi_E}{dt}$$

but, this is equal to  $i$

$$i = \epsilon_0 \frac{d\Phi_E}{dt}$$

So, an electric flux that changes with time is "equivalent" to a current

Maxwell proposed the following modified version of the 10 Ampere's law

$$\int_{\text{loop } \Gamma} \vec{B} \cdot d\vec{s} = \mu_0 i + \underbrace{\mu_0 \epsilon_0 \frac{d\Phi_E}{dt}}_{\text{called "displacement current" } i_d}$$

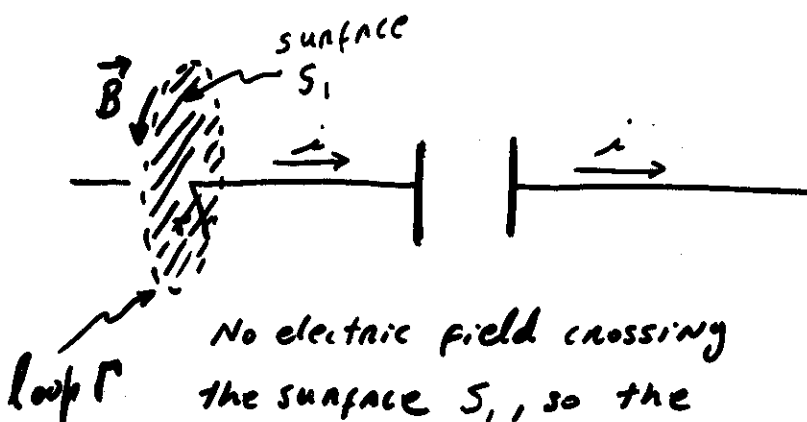
or, written more explicitly

$$\int_{\text{loop } \Gamma} \vec{B} \cdot d\vec{l} = \mu_0 i + \mu_0 \epsilon_0 \frac{d}{dt} \int_{\text{surface } S} \vec{E} \cdot d\vec{A}$$

current enclosed by  $\Gamma$

Fourth Maxwell equation (1873)

Going back to our example, let's apply the 4th Maxwell equation and find  $\vec{B}$  using the surface  $S_1$ ,



No electric field crossing the surface  $S_1$ , so the 4th Maxwell equation takes the form:

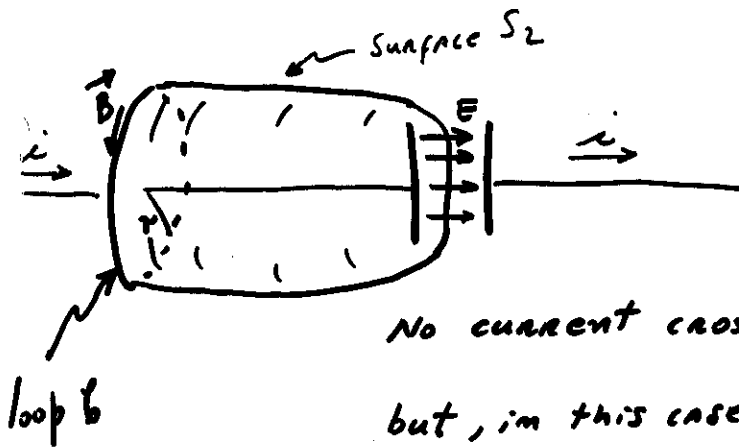
$$\int_{\text{loop } \Gamma} \vec{B} \cdot d\vec{l} = \mu_0 i$$

From which we obtain

$$B = \frac{\mu_0 i}{2\pi r}$$

③

What if we choose the surface  $S_2$ ?



No current crossing the surface  $S_2$ ,  
but, in this case, there does exist a  
electric field  $\vec{E}$  crossing the surface.  
So, the 4th Maxwell equation takes  
the form:

$$\int_{\text{loop } \Gamma} \vec{B} \cdot d\vec{l} = \mu_0 \epsilon_0 \frac{d}{dt} \int_{\text{surface } S_2} \vec{E} \cdot d\vec{A}$$
$$= EA = \frac{Q}{A\epsilon_0} A = \frac{Q}{\epsilon_0}$$

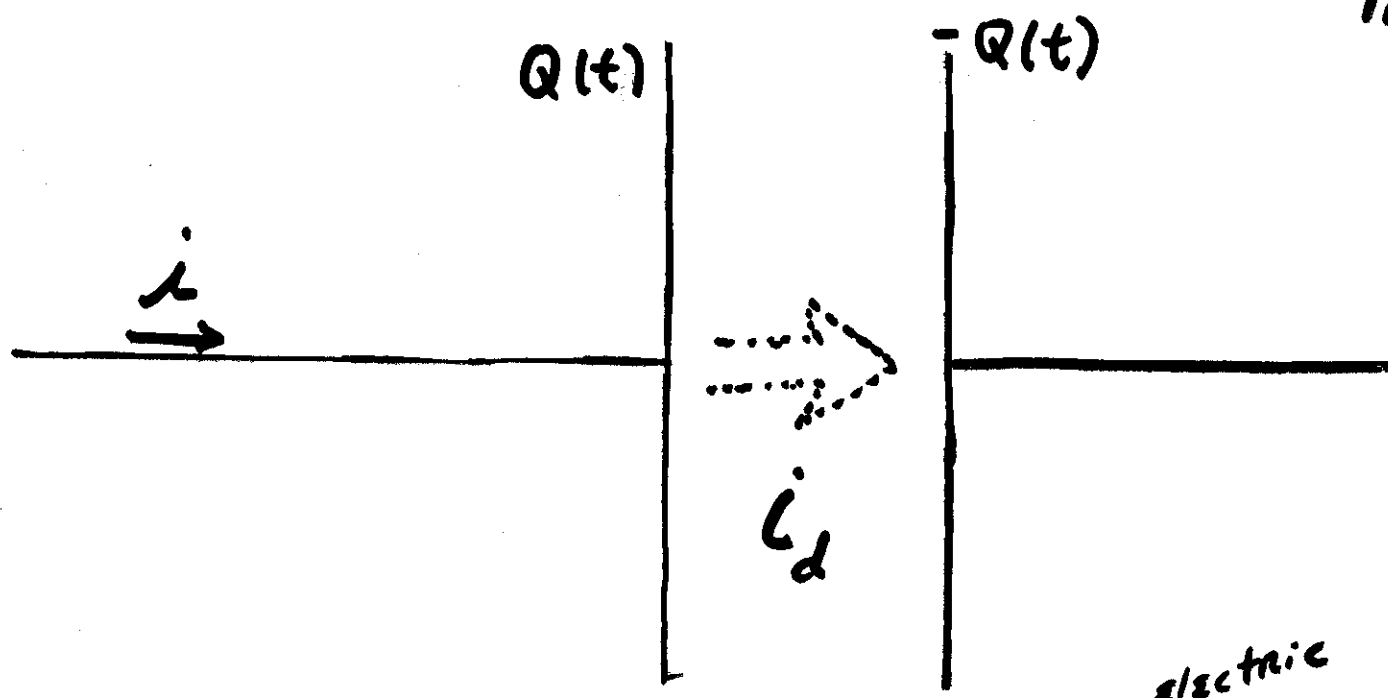
$$= \mu_0 \epsilon_0 \frac{d}{dt} \left( \frac{Q}{\epsilon_0} \right) = \mu_0 \frac{dQ}{dt} = \mu_0 i$$

$$\int_{\text{loop } \Gamma} \vec{B} \cdot d\vec{l} = \mu_0 i$$

From which we obtain

$$B = \frac{\mu_0}{2\pi r} i$$

Same result as (4)  
when the surface  
 $S_1$  was used.



charge on the plate

$$i = \frac{dQ}{dt}$$

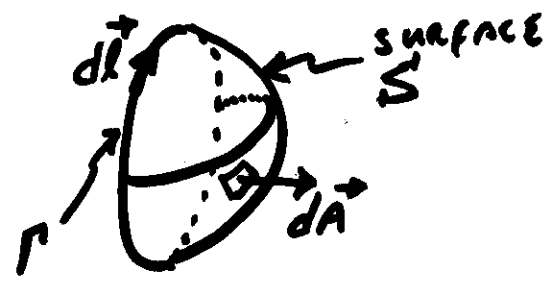
electric flux

$$i_d = \epsilon_0 \frac{d\Phi_E}{dt}$$

"displacement current"

More general,

$$\int_{\Gamma} \vec{B} \cdot d\vec{l} = \mu_0 \int_S \vec{j} \cdot d\vec{A} + \mu_0 \epsilon_0 \frac{d}{dt} \int_S \vec{E} \cdot d\vec{A}$$



4th MAXWELL'S equation