

## Relationship between the amplitudes $E_m$ and $B_m$ 13

Let's consider a sinusoidal wave

$$E = E_m \cos(kx - \omega t)$$

$$B = B_m \cos(kx - \omega t)$$

and use equation ① (see page #8 in these notes)

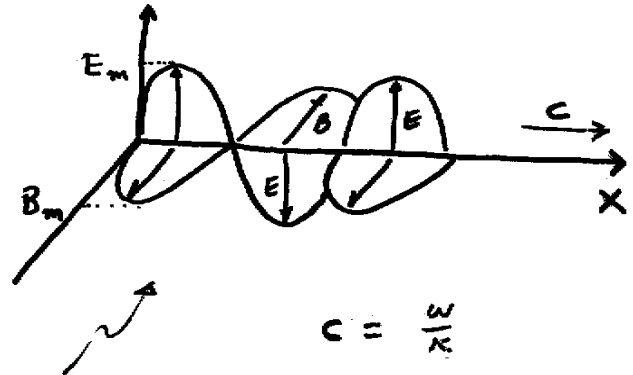
$$\frac{\partial E}{\partial x} = -k E_m \sin(kx - \omega t)$$

$$\frac{\partial B}{\partial t} = +\omega B_m \sin(kx - \omega t)$$

Eq ① implies  $\rightarrow k E_m = \omega B_m$  or

$$B_m = \frac{E_m}{c}$$

$$\frac{\partial E}{\partial x} = -\frac{\partial B}{\partial t}$$



EM propagating in the free space

$\vec{E}$  and  $\vec{B}$  are perpendicular to one another

$\vec{E} \times \vec{B}$  indicates vector product  $\rightarrow$  points in the direction where the wave propagates toward to.

## ELECTROMAGNETIC WAVES TRANSPORT ENERGY POYNTING VECTOR

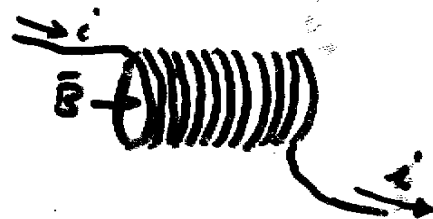
- Wherever there is an electric field  $\vec{E}$ , there is an energy density  $u_E$  (energy per unit volume, Joule/m<sup>3</sup>) associated to it:

$$u_E = \frac{1}{2} \epsilon_0 E^2 \quad (\text{Joule/m}^3)$$



- Wherever there is a magnetic field  $\vec{B}$ , there is an energy density  $u_B$  (Joules/m<sup>3</sup>) associated to it:

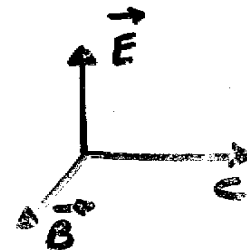
$$u_B = \frac{1}{2\mu_0} B^2$$



- In an electromagnetic wave, the electromagnetic energy density will be given by

$$u = \frac{1}{2} \epsilon_0 E^2 + \frac{1}{2\mu_0} B^2$$

we can use  $B = \frac{E}{c}$



$$= \frac{1}{2} \epsilon_0 E^2 + \frac{1}{2} \epsilon_0 E^2$$

where we have used  $c^2 = \frac{1}{\epsilon_0 \mu_0}$

else electromagnetic energy density (Joule/m<sup>3</sup>)

$$u = \epsilon_0 E^2$$

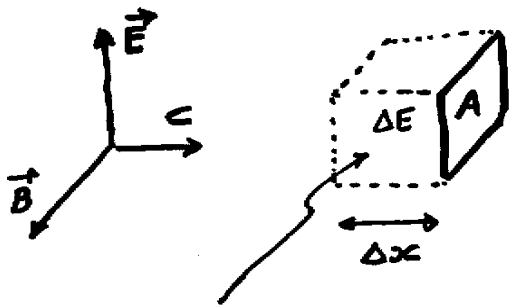
or

$$u = \frac{B^2}{\mu_0}$$

or

$$u = \frac{EB}{\mu_0 c}$$

But this energy is propagating with velocity  $c$   
Let's consider a section of area  $A$



the amount of energy in this cubic volume is  $\Delta E = \mu A \Delta x$

In an interval of time  $\Delta t = \frac{\Delta x}{c}$  all the energy in that cubic volume will have passed through the section of area  $A$

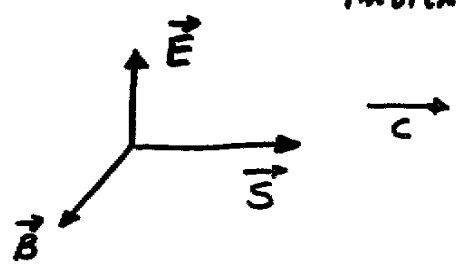
Energy per unit time crossing the area  $A$   $= \frac{\Delta E}{\Delta t} = \frac{\mu A \Delta x}{\Delta x / c} = \mu A c$

ENERGY per unit time and per unit area  $S = \mu c = \frac{E B}{\mu_0}$

$$\vec{S} = \frac{1}{\mu_0} \vec{E} \times \vec{B}$$

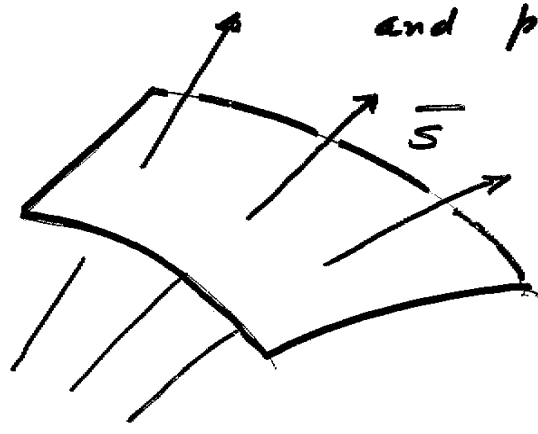
Poynting vector

indicates vector product



S is the flow of electromagnetic energy traveling through the space, per unit area and per unit time.

$$\left[ \frac{J}{m^2 sec} \right]$$



$$S = \mu c$$

$\downarrow$  energy density       $\uparrow$  speed of light

$$S = \frac{1}{\mu_0} EB$$

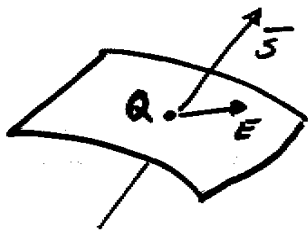
$$S = \epsilon_0 c E^2$$

$$S = \frac{c}{\mu_0} B^2$$

S: power per unit area (Watt/m<sup>2</sup>)

## AVERAGE ENERGY $\langle S \rangle$ ; INTENSITY

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$$S = \epsilon_0 c E^2$$

Let's assume the electric field at the position "Q" is of the form

$$E = E_m \cos(\omega t)$$

The Poynting vector takes the form:

$$S = \epsilon_0 c E_m^2 \cos^2(\omega t)$$



If we take average over time:

$$\langle S \rangle = \epsilon_0 c E_m^2 \underbrace{\langle \cos^2(\omega t) \rangle}_{= \frac{1}{2}}$$

$$\boxed{I = \langle S \rangle = \frac{1}{2} \epsilon_0 c E_m^2}$$

Intensity

power per  
unit area  
(Watt/m<sup>2</sup>)

Sometimes, this result is given in terms of rms value of the electric field amplitude

$$E_{rms} \equiv \frac{E_m}{\sqrt{2}} \text{ (definition)}$$

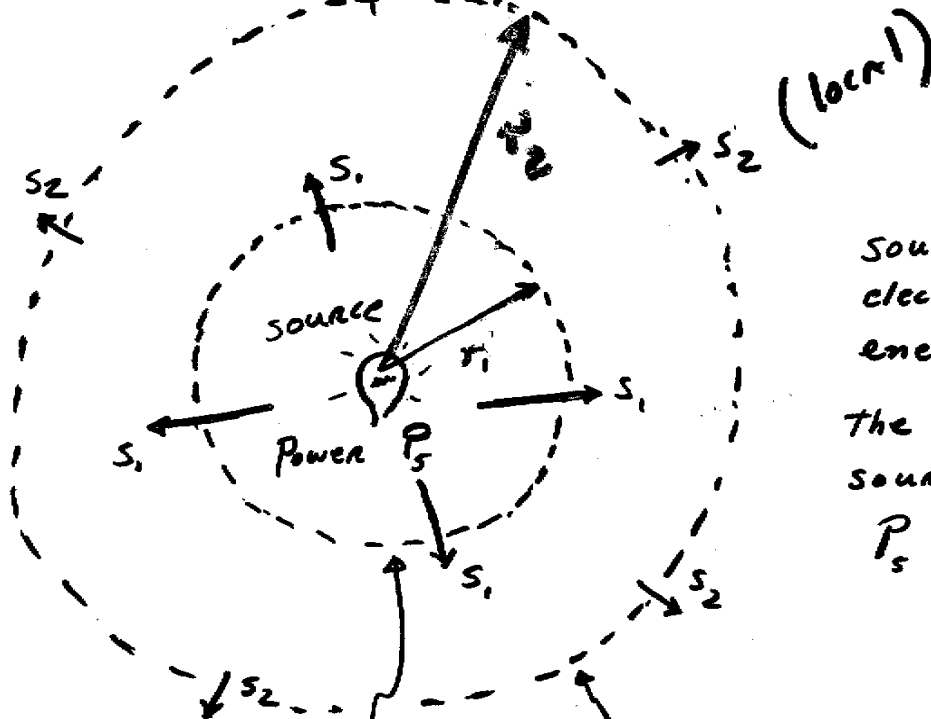
$$\boxed{I = \epsilon_0 c E_{rms}^2}$$

Intensity

(Watt/m<sup>2</sup>)

Intensity decreases with distance from the source

CASE: PUNCTUAL SOURCE emitting energy  
All over the space



Source emits  
electromagnetic  
energy

The power of the  
source is:

$$P_s \text{ (Joules/sec)}$$

Notice, conservation of energy implies:

$$P_s = \langle S_1 \rangle 4\pi r_1^2 = \langle S_2 \rangle 4\pi r_2^2$$

$$= I_1 4\pi r_1^2 = I_2 4\pi r_2^2$$

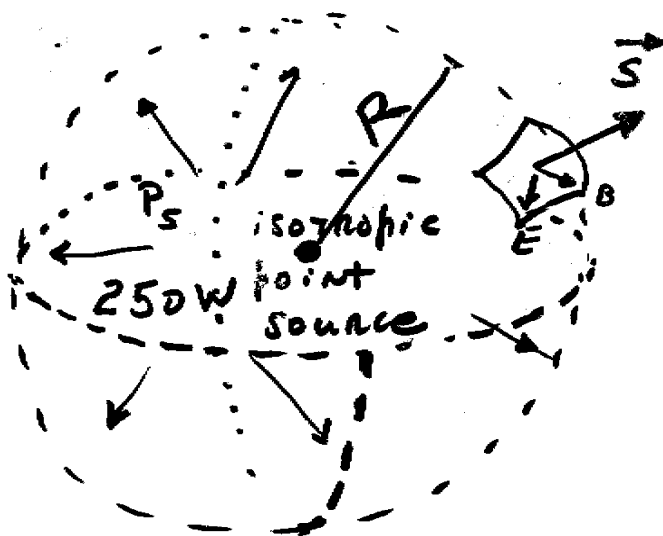
That is, for a constant value of  $P_s$ , the intensity  
varies with the distance  $r$ :

$$I(r) = \frac{P_s}{4\pi r^2}$$

Remember:  
 $I = \langle S \rangle$

Example: sample problem 34-1 (pg 811)

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$$\left[ \frac{\text{Joule}}{\text{m}^2 \text{ sec}} \right]$$

$$\vec{S} = \frac{1}{\mu_0} \vec{E} \times \vec{B}$$

$$E = E_m \cos(\omega t)$$

$$\langle S \rangle = \frac{1}{2} \epsilon_0 c E_m^2$$

$$\frac{250 \text{ W}}{4\pi R^2} = \langle S \rangle = \frac{1}{2} \epsilon_0 c E_m^2 \quad *$$

$$R = 1.8 \text{ m}$$

$$\epsilon_0 =$$

$$c =$$

$$\Rightarrow E_m^2 = \frac{2 \times 250 \frac{\text{J}}{\text{s}}}{4\pi (1.8 \text{ m})^2 \epsilon_0 c}$$

Electric field  $E = \frac{\text{Volts}}{\text{m}}$  (Remember  $E \cdot d = V$ )

$\text{Volt} = \frac{\text{work}}{\text{unit charge}} = \frac{J}{C}$

$\epsilon_0 = 8.85 \times 10^{-12} \frac{F}{m}$   
 $= 8.85 \times 10^{-12} \frac{C^2}{Jm}$

Remember  
 $F = \frac{1}{4\pi\epsilon_0} \frac{q^2}{r^2}$   
So,  
 $\epsilon_0 = \frac{C^2}{Nm^2} = \frac{C^2}{Jm}$

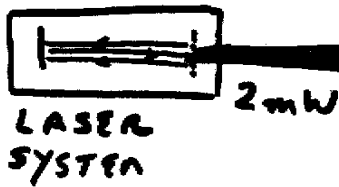
Therefore

$E_m^2 = \frac{2 \times 250 \frac{J}{s}}{4\pi (1.8m)^2 \cdot 8.85 \cdot 10^{-12} \frac{C^2}{Jm} \times 3 \times 10^8 \frac{m}{s}}$   
 $= 0.46 \times 10^4 \frac{J^2}{m^2 C^2} = 0.46 \times 10^4 \frac{V^2}{m^2}$

$E_m = 0.68 \times 10^2 \frac{V}{m}$

In term of rms:  $E_{rms} = \frac{E_m}{\sqrt{2}} = 0.48 \frac{V}{m} \times 10^2$

\* Exercise: The beam diameter of a typical <sup>231</sup> He-Ne laser is  $\sim 2 \text{ mm}$ .  
 (just at the output of the laser system case)



Find the amplitudes of the electric field and magnetic field of the laser beam.

$$I = \langle S \rangle = \frac{1}{2} \epsilon_0 c E_m^2 = \frac{\text{Power}}{\text{unit area}} = \frac{2 \text{ mW}}{\pi (1 \text{ mm})^2}$$

$$= \frac{2 \times 10^{-3} \text{ J/s}}{\pi (10^{-3})^2 \text{ m}^2} = 0.64 \times 10^3 \frac{\text{J}}{\text{sm}^2}$$

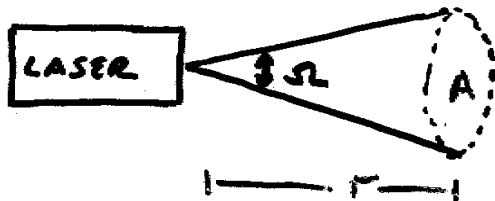
$$\Rightarrow E_m^2 = \frac{2 \times 0.64 \times 10^3 \frac{\text{J}}{\text{sm}^2}}{\epsilon_0 c}$$

$$E_m = \sqrt{48.2 \times 10^4} \text{ V}$$

$$E_m = 7 \times 10^2 \frac{\text{V}}{\text{m}}$$

Note: Output laser beam is not perfectly collimated.

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Solid angle  $\Omega$   
is measured in radians

$$\Omega = \frac{A}{r^2}$$

where

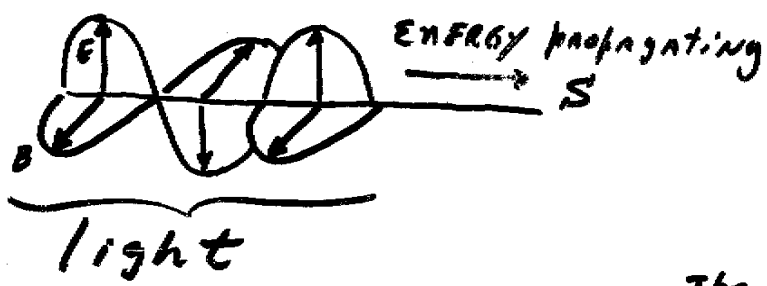
A: cross section area

r: distance

# LIGHT as a WAVE



- Light is a wave of electric fields and magnetic fields
- The vector product  $\vec{E} \times \vec{B}$  gives the direction of light propagation



$$\langle S \rangle = \frac{1}{2} \epsilon_0 c E_m^2$$

CLASSICAL VIEW



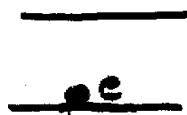
The higher  $E_m$   
the higher the power  
propagating in  
the wave.

However, all this nice theory could not explain the following experimental result:

■ Small amplitude  $E_m$



ATOM



higher  $E_m$



no matter how much you increase the incident light power (increase of  $E_m$ ) the electron is not excited to the upper level of energy.

■ HOWEVER

§ ~~~~~

higher frequency light

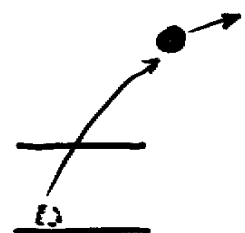
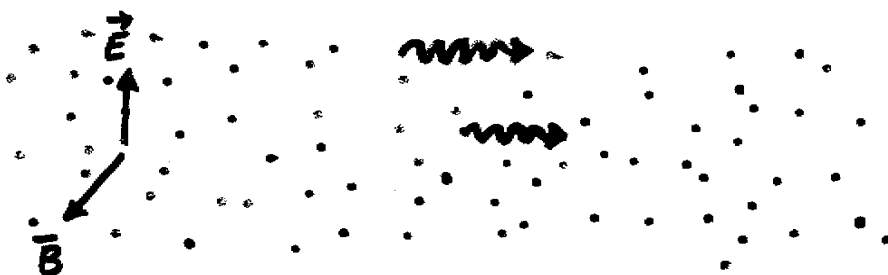


Photo-electric effect

how to explain it ?

# LIGHT as a Radiation Field of Quantized energy packets



Light is made out of photons.

$$1 \text{ photon energy} = h f = \frac{hc}{\lambda}$$

$\uparrow$                        $\uparrow$   
 frequency of  
 the radiation

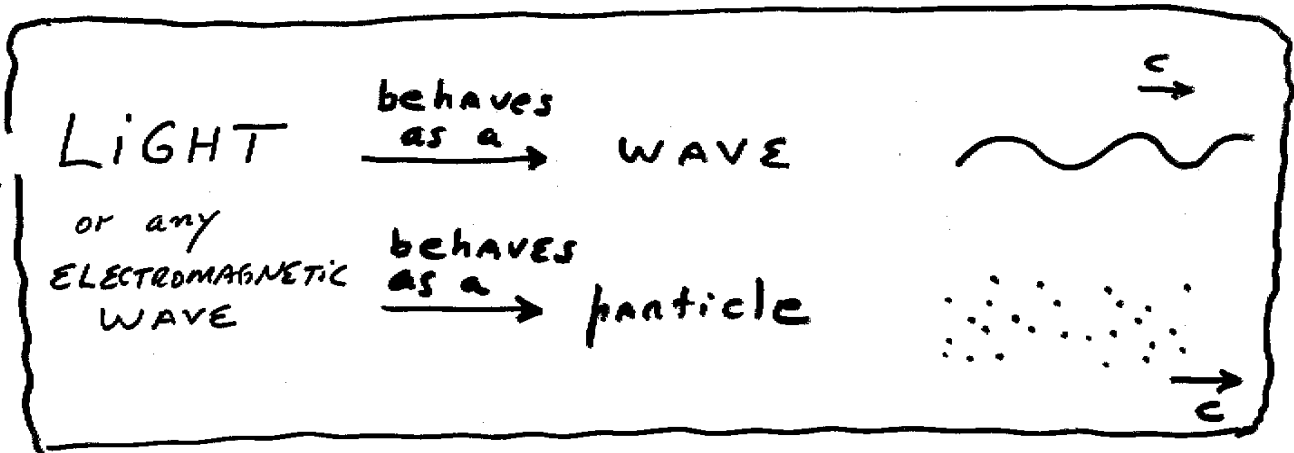
constant #

"Planck constant"

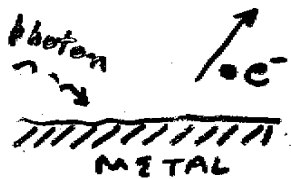
$$= 6.6 \times 10^{-34} \text{ Joule} \cdot \text{sec}$$

$$= 4.14 \times 10^{-15} \text{ eV} \cdot \text{sec}$$

\* Exercise: Find the energy of 1 photon whose wavelength is  $\lambda = 500 \text{ nm}$  ("green light")



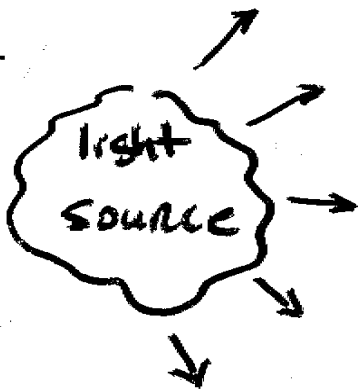
The particle nature of light was first proposed by Albert Einstein in 1905 in his explanation of the "photoelectric effect."



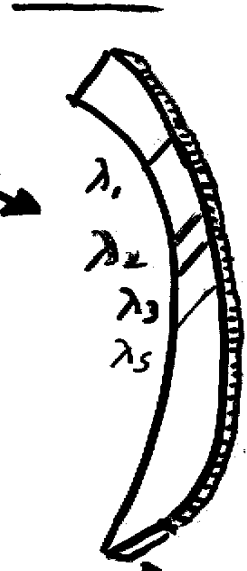
$$\begin{aligned}
 E_{\text{photon}} &= hf = 6.6 \times 10^{-34} \text{ Joule} \cdot \text{sec} \cdot f \\
 &= 4.14 \times 10^{-15} \text{ eV} \cdot \text{sec} \cdot f
 \end{aligned}$$

$$= h \frac{c}{\lambda} = \frac{1240 \text{ eV} \cdot \text{nm}}{\lambda}$$

$$\begin{aligned}
 \text{nm} &= \text{nano-meter} \\
 &= 10^{-9} \text{ m}
 \end{aligned}$$



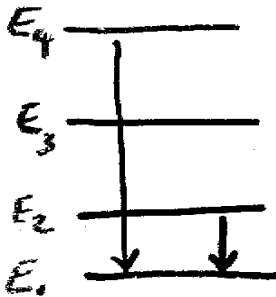
Experimental observation of a discrete lines spectra



photographic film

Hypothesis

$$\Delta E = \frac{hc}{\lambda}$$



ATOM



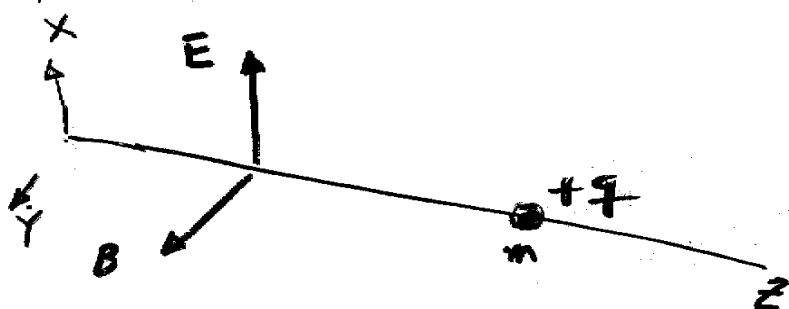
Measurements of discrete values of  $\lambda$  allowed to make models of the ATOM

## LIGHT WAVE-particle duality

the propagation of light is governed by its wave properties

the light-matter interaction is governed by its particle-like character.

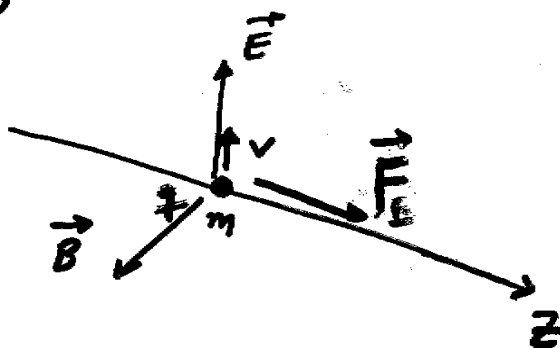
If light behaves as a particle,  
does it have  
momentum ?  
angular momentum ?



$\vec{E}$   $\longrightarrow$  will shake the charge up and down  
(along the x-axis)

What about  $\vec{B}$  ?

$$\vec{F}_B = q \vec{v} \times \vec{B}$$



$\vec{F}_B$  acts in the  
direction of the propa-  
gation of light

So, the mass  $m$  will  
gain some linear  
momentum along the  
z-axis

$$F = q v B$$

$$B = \frac{E}{c}$$

$$= \frac{q v}{c} E$$

velocity  $\times$  force  
= power

= work  $W$  done (by the electric field  $E$ ) on the charge  $q$ , per unit time

$$F = \frac{1}{c} \frac{dW}{dt} = \frac{d}{dt} \left( \frac{W}{c} \right)$$

$W$  must come from the energy of the light

momentum

if one photon is absorbed

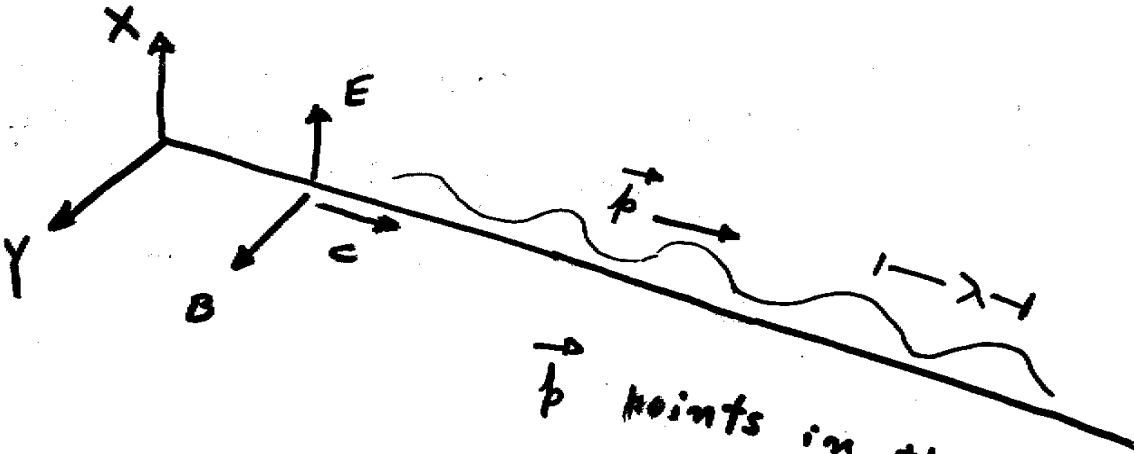
$$p = \frac{W}{c} = \frac{h f}{c}$$

$$W = h f$$

$h$  is the Planck's constant

$$p = \frac{h}{\lambda}$$

Linear momentum carried by a photon



$\vec{p}$  points in the direction  
of the light propagation

$$p = \frac{h}{\lambda}$$

\* Exercise: Problems 13E, 14E, 16P; Chap 34, p. 828  
6th Ed

## MAXWELL'S RAINBOW

