

INTRODUCTION TO QUANTUM MECHANICS

Homework - 5

Due: 11-12-2009

1. The evolution operator $\tilde{U}(t_2, t_1)$ acting on a state $|\Psi\rangle$ at a given time t_1 gives the state at the time t_2 . Symbolically,

$$\underbrace{\tilde{U}(t_2, t_1)}_{\text{Operator}} \underbrace{|\Psi(t_1)\rangle}_{\text{State}} = \underbrace{|\Psi(t_2)\rangle}_{\text{State}}$$

Assuming that only 4 base states $\{|n\rangle; n = 1, 2, 3, 4\}$ are necessary to completely describe a quantum system:

- 1.A Construct explicitly the corresponding matrix $[U](t_2, t_1)$ representing the $\tilde{U}(t_2, t_1)$ operator.

- 1.B Find the form that the matrix $[U](t_2, t_1)$ adopts when $t_2 \rightarrow t_1$
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2. Assume that only 4 base states $\{|n\rangle; n = 1, 2, 3, 4\}$ are necessary to completely describe a quantum system. That is, a general state $|\Psi(t)\rangle$ can

be expressed as $|\Psi(t)\rangle = \sum_{n=1}^4 |n\rangle A_n(t)$.

We know that, the value of a particular amplitude $A_m(t + \Delta t)$ at the time $t + \Delta t$ can be calculated in terms of the four amplitudes $A_n(t)$ ($n = 1, 2, 3$ and 4) evaluated at the previous time t . Indeed, for small Δt we have,

$$A_m(t + \Delta t) = \sum_{n=1}^4 \left(\delta_{mn} - \frac{i}{\hbar} (\Delta t) H_{mn}(t) \right) A_n(t)$$

- 2.A For each value of m (i.e. 1, 2, 3, and 4) evaluate $\sum_{n=1}^4 \delta_{mn} A_n(t)$

Subsequently, express the result in a matrix format

2.B For each value of m (i.e. 1, 2, 3, and 4) evaluate $\sum_{n=1}^4 \left(\frac{i}{\hbar} \Delta t H_{mn}(t) \right) A_n(t)$

Subsequently, express the result in a matrix format

3. In the ammonia molecule, introduced in section 7.1, we identify two possible states:

State $|1\rangle$ the “down” state (corresponding to a charge distribution when the nitrogen is *below* the plane defined by the three hydrogen atoms.)

State $|2\rangle$ the “up” state (corresponding to a charge distribution when the nitrogen is *above* the plane defined by the three hydrogen atoms.)

The evolution in time of a wavefunction

$$|\psi(t)\rangle = |1\rangle A_1(t) + |2\rangle A_2(t) \quad (1)$$

is governed by

$$i\hbar \frac{dA_1(t)}{dt} = E_0 A_1 - W A_2$$

$$i\hbar \frac{dA_2(t)}{dt} = -W A_1 + E_0 A_2$$

Let's define the “new” states $|S_I\rangle$ and $|S_{II}\rangle$ in terms of the “old” states $|1\rangle$ and $|2\rangle$ according to:

$$|S_I\rangle \equiv \frac{1}{\sqrt{2}} (|1\rangle + |2\rangle) \quad , \quad |S_{II}\rangle \equiv \frac{1}{\sqrt{2}} (|1\rangle - |2\rangle)$$

In terms of the new states, the wavefunction can also be expressed as

$$|\psi(t)\rangle = |S_I\rangle B_I(t) + |S_{II}\rangle B_{II}(t) \quad (2)$$

3A Evaluate $\langle S_I | S_I \rangle$, $\langle S_{II} | S_{II} \rangle$, $\langle S_I | S_{II} \rangle$

(Assume the states $|1\rangle$ and $|2\rangle$ are ortho-normalized.)

3B Express each of the two B -amplitudes in terms of the A -amplitudes

3C Find explicitly $B_I(t)$ and $B_{II}(t)$ as a function of t , and in terms of E_0 and W .

3D Write down the matrix representation of the Hamiltonian operator when using the base $\{|1\rangle, |2\rangle\}$

Write down the matrix representation of the same Hamiltonian operator but when using the base $\{|S_I\rangle, |S_{II}\rangle\}$

4. Dynamics of an electron in a one-dimensional crystal

By assuming that the electron will jump only from one atom to the neighbor of either side (with an amplitude probability equal to $-W$), one arrives to the following Hamiltonian equations.

$$i\hbar \frac{dA_n}{dt} = -W A_{n-1} + E_o A_n - W A_{n+1} \quad (3)$$

4.1 What is the meaning (or interpretation) of E_o in these equations

4.2 Show that these equations have solutions of the form

$$|\phi^{(E)}(t)\rangle = \sum_n |n\rangle A_n^{(E)}(t) = \sum_n |n\rangle a_n e^{-(i/\hbar)(E)t} \quad (4)$$

What is the meaning of E in this solution (4)

What is the condition to be satisfied by a_n (for (4) to become a solution of Eq. (3))?

What is the physical meaning of a_n in the solution (4)

5. Dynamics of an electron in a one-dimensional crystal

For the case of a low energy wave-packet whose main wavenumber k is around k_o , show explicitly that

$$m_{eff} v_g = \hbar k_o$$

(in the jargon of solid state physics, $\hbar k$ is called the crystal momentum.)

6. Dynamics of an electron in a one-dimensional crystal

For a free particle $E(k) = \frac{1}{2m} \hbar^2 k^2$. From this expression, the mass m can be obtained by differentiating $E(k)$ twice. That is,

$$m = \frac{\hbar^2}{\frac{d^2}{dk^2} E(k)}$$

Consider now the case of an electron in a lattice, where

$$E(k) = 2W - 2W \cos(kb)$$

- 6A** Use the definition given above to obtain the effective mass m^* associated to an electron in a lattice. That is, evaluate

$$m^* = \frac{\hbar^2}{\frac{d^2}{dk^2} E(k)}$$

- 6B** Evaluate the group velocity of a wavepacket describing the motion of an electron in a lattice
- 6C** Evaluate the product of the effective mass with the group velocity.

Verify that for small values of k (compared to π/b) the product

$$m^* v_g \xrightarrow{\text{small } k} \hbar k$$